

# Modified Block BiCGSTAB for Lattice QCD

Yoshifumi Nakamura

RIKEN AICS

10 Jun 2011



# Introduction

**Motivation:** Understanding particle physics

# Introduction

**Motivation:** Understanding particle physics

Non-perturbative approach of Quantum Chromo Dynamics(QCD)

# Introduction

**Motivation:** Understanding particle physics

Non-perturbative approach of Quantum Chromo Dynamics(QCD)

Lattice QCD (LQCD)

# Introduction

**Motivation:** Understanding particle physics

Non-perturbative approach of Quantum Chromo Dynamics(QCD)

Lattice QCD (LQCD)

But huge CPU resources require for LQCD simulations

# Introduction

**Motivation:** Understanding particle physics

Non-perturbative approach of Quantum Chromo Dynamics(QCD)

Lattice QCD (LQCD)

But huge CPU resources require for LQCD simulations

It is important **to reduce cost for LQCD simulations**

# Contents

- QCD and LQCD
- Modified Block BiCGSTAB for Lattice QCD

# Quantum Chromo Dynamics(QCD)

- basic theorem of hadron physics
- describing interaction for **quark** and **gluon**
- typical scale is **0.000 000 000 001 m = 1 fm**

hadron(color-neutral): meson( $\pi, K \dots$ ), baryon(p,n...)



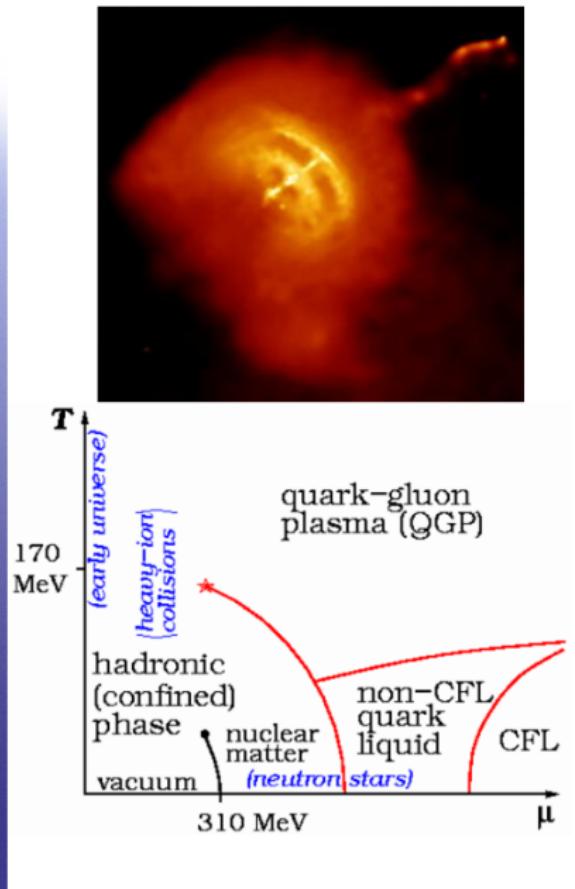
quark: 3 (R, G, B)

gluon: 8 (mass 0, charge 0, spin 1)

# quark: 6 flavours

	u(up)	c(charm)	t(top)
mass	1.7-3.3 MeV	$1.27^{+0.07}_{-0.09}$ GeV	172.0(22) GeV
charge	2/3	2/3	2/3
spin	1/2	1/2	1/2
	d(down)	s(strange)	b(bottm)
mass	4.1-5.8 MeV	$101^{+29}_{-21}$ MeV	$4.19^{+0.18}_{-0.06}$ GeV
charge	-1/3	-1/3	-1/3
spin	1/2	1/2	1/2

# History of the Universe



images from NASA and wikipedia

# QCD action

$$L = \sum_i \bar{\psi}_i D(\mathbf{m}_i) \psi_i - \frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu}$$

$$\begin{aligned} F_{\mu\nu}^a &: \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + i \mathbf{g} f^{abc} A_\mu^b A_\nu^c \\ D(m_i) &: \gamma_\mu (i \partial_\mu + \mathbf{g} A_\mu^a T^a) - \mathbf{m}_i \end{aligned}$$

# Path integral

Expectation value of observable  $O$ :

$$\langle O \rangle = \frac{1}{Z} \int dA d\bar{\psi} d\psi \ e^{i \int d^4x L(x,t)} \ O(A, \bar{\psi}, \psi)$$

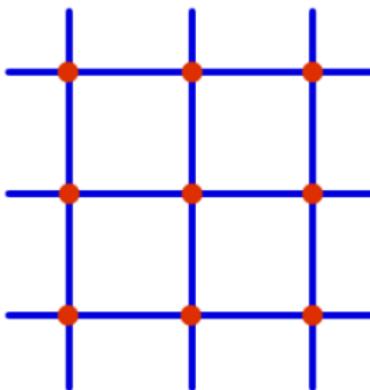
QED: perturbation + renormalization

QCD: **perturbation does not work** at low energy since coupling constant is big

# Lattice QCD

Non-perturbative approach to solving QCD

- space - time discretization
  - quark field: color  $\times$  spinor / site  $\rightarrow$  12 complex numbers
  - gluon field: SU(3) matrix / link  $\rightarrow$  9 complex numbers



# Lattice QCD simulation

fermion field(Grassmann number)

→ pseudo-fermion field(usual number)

$$\int d\bar{\psi}_i d\psi_i \exp(-\sum_{i=1}^2 \bar{\psi}_i \mathbf{D}(\mathbf{m}_i) \psi_i) = \det \mathbf{D}(\mathbf{m}_1) \det \mathbf{D}(\mathbf{m}_2)$$

when  $\mathbf{m} = \mathbf{m}_1 = \mathbf{m}_2$

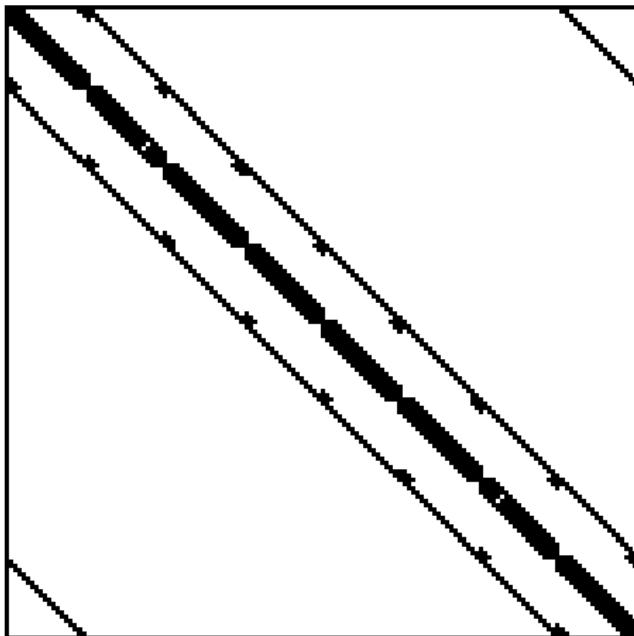
$$\det \mathbf{D}(\mathbf{m})^2 = \int d\phi_i^\dagger d\phi_i \exp(-\sum_i \phi_i^\dagger [\mathbf{D}^\dagger(\mathbf{m}) \mathbf{D}(\mathbf{m})]^{-1} \phi_i)$$

condition number increases as  $\mathbf{m}$  decreasing

$\mathbf{D}$  is  $12V \times 12V$  complex sparse matrix

e.g.  $V = 32^3 \times 64 \Rightarrow O(10^7)$

# Sparsity pattern of Wilson-Dirac matrix



$$n = 8^4 \times 12$$

# Modified Block BiCGSTAB for Lattice QCD

Y. Nakamura, K. -I. Ishikawa, Y. Kuramashi, T. Sakurai, H. Tadano (2011)

## Outline

- Krylov subspace method
- Block Krylov subspace method
- Algorithm of Modified Block BiCGSTAB for Lattice QCD
- Numerical test results
- Summary

# Krylov subspace method

iterative method to solve system of linear equations

$$Ax = b$$

$$x = A^{-1}b$$

$$\begin{bmatrix} a_{1,1} & \cdots & a_{1,n} \\ \vdots & \ddots & \vdots \\ a_{n,1} & \cdots & a_{n,n} \end{bmatrix} \begin{bmatrix} x_1 \\ \vdots \\ x_n \end{bmatrix} = \begin{bmatrix} b_1 \\ \vdots \\ b_n \end{bmatrix}$$

by using matrix-vector multiplication

# Krylov subspace

$$\mathcal{K}_k \equiv \text{span}(v, Av, A^2v, \dots, A^{k-1}v)$$

- 1 guess initial approx. solution vector  $x_0$  for  $Ax = b$
- 2 renew approx. solutions  $x_1, x_2, x_3, \dots$  with keeping condition of  $x_k - x_0 \in \mathcal{K}_k(A, r_0) = \text{span}(r_0, Ar_0, \dots, A^{k-1}r_0)$

First residual vector:

$$r_0 = b - Ax_0$$

Approximate solution:

$$x_k = x_0 + \sum_{i=0}^{k-1} c_i A^i r_0$$

# Conjugate gradient (CG) method

Hestenes, Stiefel (1952)

Krylov subspace method for symmetric positive definite

- minimize  $f(x) = (x, Ax) - 2(x, b)$
- (k+1)-th approx. solution vector:  $x_{k+1} = x_k + \alpha p_k$
- (k+1)-th research vector:  $p_{k+1} = r_{k+1} + \beta p_k$
- $(p_{k+1}, Ap_k) = 0$

(k+1)-th residual vector:

$$r_{k+1} = r_k - \alpha A p_k$$

# Properties of CG method

- conjugate property

$$(p_i, A p_j) = 0, \quad i \neq j$$

- orthogonality

$$(r_i, r_j) = 0, \quad i \neq j$$

System converges theoretically by 'n' iteration at most

# Block Krylov subspace method

$$AX = B$$

$$X = [x^{(1)}, x^{(2)}, \dots, x^{(L)}]$$

$$B = [b^{(1)}, b^{(2)}, \dots, b^{(L)}]$$

$$X_k - X_0 \in \mathcal{K}_k(A, R_0) = \text{span}(R_0, AR_0, \dots A^{k-1}R_0)$$

Approx. solutions  $X_k$ :

$$X_k = X_0 + \sum_{i=0}^{k-1} A^i R_0 \gamma_i$$

$\gamma_i$  is  $L \times L$  matrix

To solve  $x^{(i)}$ , one can use information of  $\mathcal{K}_k(A, r_0^{(j)})$

# Better cache usage

at Matrix Vector multiplication (MVM)

$$\begin{bmatrix} w_1^{(1)} & w_1^{(2)} \\ \vdots & \vdots \\ w_n^{(1)} & w_n^{(2)} \end{bmatrix} = \begin{bmatrix} a_{1,1} & \cdots & a_{1,n} \\ \vdots & \ddots & \vdots \\ a_{n,1} & \cdots & a_{n,n} \end{bmatrix} \begin{bmatrix} v_1^{(1)} & v_1^{(2)} \\ \vdots & \vdots \\ v_n^{(1)} & v_n^{(2)} \end{bmatrix}$$

One can calculate  $a_{1,1} \times v_1^{(2)}$  (not  $a_{1,2} \times v_2^{(1)}$ ) right after  $a_{1,1} \times v_1^{(1)}$

# Wilson-Dirac operator

$$A\phi = \sum_{x=1}^{L_x \times L_y \times L_z \times L_t} (\phi_x - \kappa \eta_x), \quad \eta_x = \sum_{\mu=1}^4 [(1 - \gamma_\mu) U_{x,\hat{\mu}} \phi_{x+\hat{\mu}} + (1 + \gamma_\mu) U_{x-\hat{\mu},\hat{\mu}}^\dagger \phi_{x-\hat{\mu}}]$$

To computer  $\eta_x$  (hopping term multiplication)

Flops : 1320

store : 12 complex numbers

load : 72 + 96 complex numbers, for  $U$  and  $\phi$

# Wilson-Dirac operator

$$A\phi = \sum_{x=1}^{L_x \times L_y \times L_z \times L_t} (\phi_x - \kappa \eta_x), \quad \eta_x = \sum_{\mu=1}^4 [(1 - \gamma_\mu) U_{x,\hat{\mu}} \phi_{x+\hat{\mu}} + (1 + \gamma_\mu) U_{x-\mu,\hat{\mu}}^\dagger \phi_{x-\hat{\mu}}]$$

To computer  $\eta_x$  (hopping term multiplication)

Flops : 1320

store : 12 complex numbers

load : 72 + 96 complex numbers, for  $U$  and  $\phi$

Hopping term mult. for multiple right hand sides

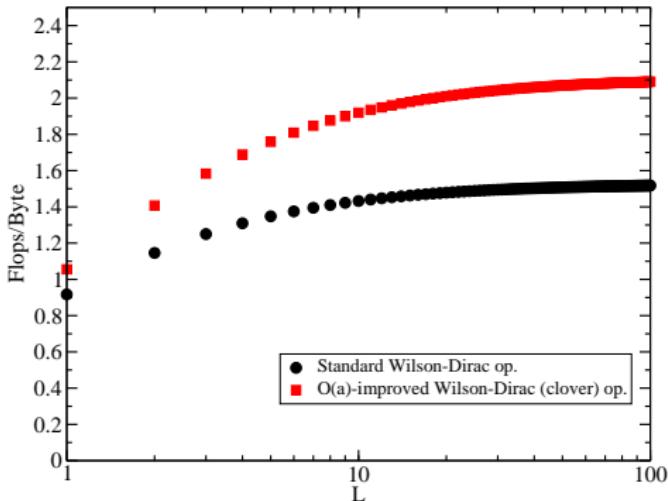
$$\eta_x^{(1, \dots, L)} = \sum_{\mu=1}^4 [(1 - \gamma_\mu) \textcolor{red}{U}_{x,\hat{\mu}} \phi_{x+\hat{\mu}}^{(1, \dots, L)} + (1 + \gamma_\mu) \textcolor{red}{U}_{x-\mu,\hat{\mu}}^\dagger \phi_{x-\hat{\mu}}^{(1, \dots, L)}]$$

Size of 8  $U$  is 576 (1152) bytes in the single (double) precision

Able to keep in low level cache and use  $L$  times

# Dirac op. Flops/Byte

with the single precision for multiple right hand sides



Block Krylov subspace method is suitable for recent high performance computer architecture

# Numerical difficulty

Dirac matrix in lattice QCD is non-Hermitian

# Numerical difficulty

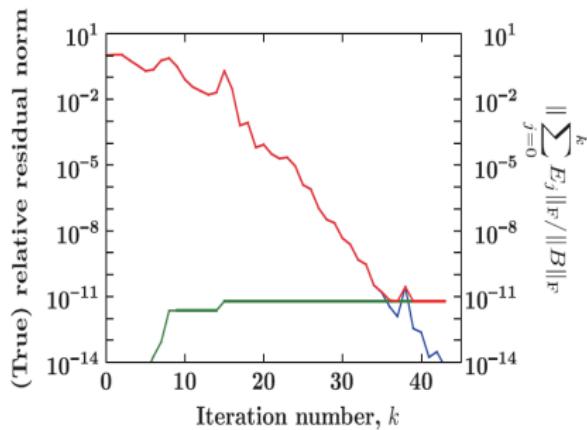
Dirac matrix in lattice QCD is non-Hermitian

→ Block BiCGSTAB (A. El Guennouni, K. Jbilou, H. Sadok (2003))

# Numerical difficulty

Dirac matrix in lattice QCD is non-Hermitian

→ Block BiCGSTAB (A. El Guennouni, K. Jbilou, H. Sadok (2003))

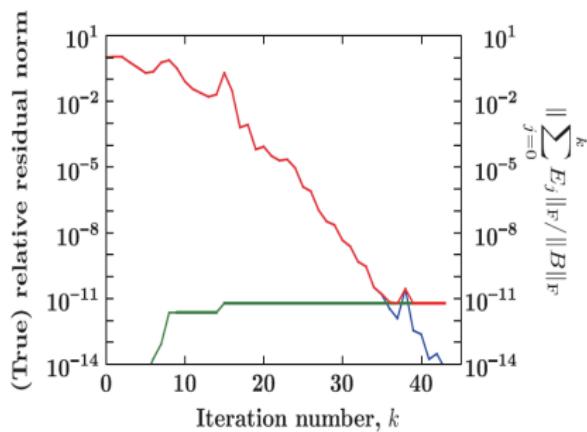


- Block BiCGSTAB has numerical error due to multiple right hand sides

# Numerical difficulty

Dirac matrix in lattice QCD is non-Hermitian

→ Block BiCGSTAB (A. El Guennouni, K. Jbilou, H. Sadok (2003))



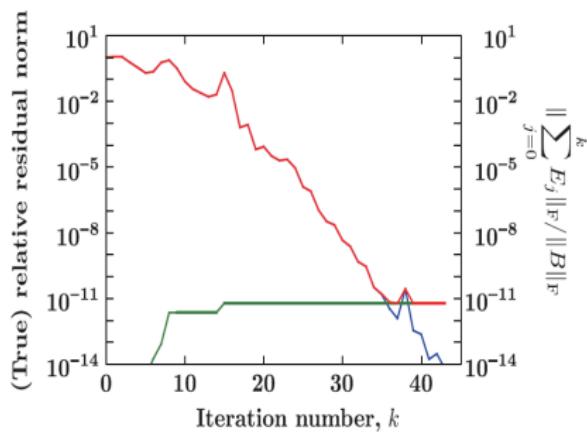
- Block BiCGSTAB has numerical error due to multiple right hand sides
- Block BiCGGR improved this problem significantly

Tadano, Sakurai, Kuramashi (2009)

# Numerical difficulty

Dirac matrix in lattice QCD is non-Hermitian

→ Block BiCGSTAB (A. El Guennouni, K. Jbilou, H. Sadok (2003))



- Block BiCGSTAB has numerical error due to multiple right hand sides
- Block BiCGGR improved this problem significantly

Tadano, Sakurai, Kuramashi (2009)

Block BiCGGR sometimes fails to converge

Further robustness and convergence are needed!!

# Modified Block BiCGSTAB algorithm

```
1  init.  $X \in \mathbb{C}^{N \times L}$ 
2   $R = B - AX$ 
3   $P = R$ 
4  choose.  $\tilde{R} \in \mathbb{C}^{N \times L}$ 
WHILE  $\max_i(|r^{(i)}|/|b^{(i)}|) \leq \epsilon$ 
4.1 QR decomp  $P = Q\gamma, P \leq Q$ 
4.2  $U = MP$ 
4.3  $V = AU$ 
4.4 solve  $(\tilde{R}^H V)\alpha = \tilde{R}^H R$  for  $\alpha$ 
4.5  $T = R - V\alpha$ 
4.6  $S = MT$ 
4.7  $Z = AS$ 
4.8  $\zeta = \text{Tr}(Z^H T_k)/\text{Tr}(Z^H Z_k)$ 
4.9  $X = X + U\alpha + \zeta S$ 
4.10  $R = T - \zeta Z$ 
4.11 solve  $(\tilde{R}^H V)\beta = -\tilde{R}^H Z$  for  $\beta$ 
4.12  $P = R + (P - \zeta V)\beta$ 
END
```

- by QR decomposition,  
numerical error ↘  
convergence ↗
- minimize comm. overhead  
by domain decomposition  
preconditioning with single  
precision acceleration

# Preconditioning

Original linear system:

$$Ax = b$$

Preconditioned system:

$$x = My$$

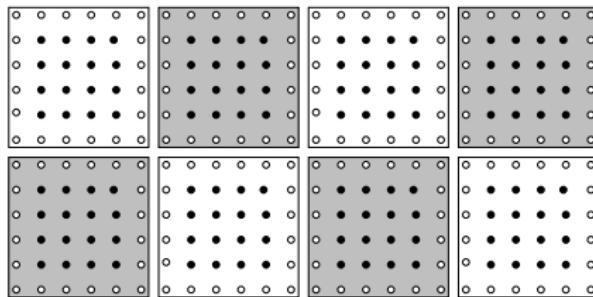
$$AMy = b$$

Preconditioner

$$M \approx A^{-1}$$

Condition number:  $AM < A$

# Domain decomposition



$$M_{SAP} = K \sum_{j=0}^{N_{SAP}} (I - AK)^j,$$

$$K = \begin{pmatrix} B_{EE} & \mathbf{0} \\ -B_{OO}A_{OE}B_{EE} & B_{OO} \end{pmatrix}$$

$B_{EE}$  ( $B_{OO}$ ) is an approximation for  $A_{EE}^{-1}$  ( $A_{OO}^{-1}$ )

# Single precision acceleration

“sloppy” precision can be used in right preconditioning

Suppose: calculation of  $S = MT$  at line 4.6 in Algorithm is performed with “sloppy” precision in  $k$ -th iteration

$$\begin{aligned} S_k &\rightarrow S'_k = S_k + \delta S_k \\ Z_k &\rightarrow Z'_k = AS'_k \\ \zeta_k &\rightarrow \zeta'_k = \zeta_k + \delta \zeta_k \\ X_{k+1} &\rightarrow X'_{k+1} = X_k + U_k \alpha_k + \zeta'_k S'_k \end{aligned}$$

These yield

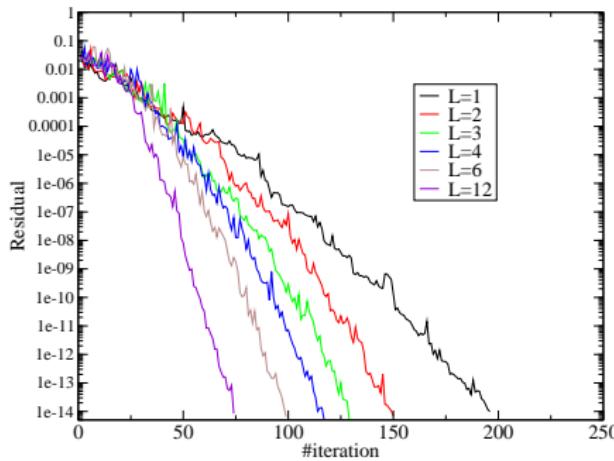
$$\begin{aligned} R'_{k+1} &= R_k - V_k \alpha_k - \zeta'_k Z'_k \\ &= R_k - AU_k \alpha_k - \zeta'_k AS'_k \\ &= B - AX_k - A(U_k \alpha_k + \zeta'_k S'_k) \\ &= B - AX'_{k+1} \end{aligned}$$

# Numerical test

lattice size :  $32^3 \times 64$   
quark masses : almost physical  
statistics : 10 independent configurations  
platform : T2K-Tsukuba 16 nodes

---

T2K-Tsukuba : quad-socket, 2.3GHz Quad-core AMD Opteron  
: 64KB/core L1\$, 512KB/core L2\$, 2MG/chip L3\$  
: 8GB DDR2-667 /socket



# Results

$L \times 12/L$	time[s]	T(gain)	NMVM	NM(gain)
<b><math>1 \times 12</math></b>	3827(755)	1	17146(3326)	1
<b><math>2 \times 6</math></b>	2066(224)	1.9	12942(1379)	1.3
<b><math>3 \times 4</math></b>	1619(129)	2.4	10652(832)	1.6
<b><math>4 \times 3</math></b>	1145(99)	3.3	9343(835)	1.8
<b><math>6 \times 2</math></b>	1040(87)	3.7	7888(663)	2.2
<b><math>12 \times 1</math></b>	705(70)	<b>5.4</b>	6106(633)	<b>2.8</b>

- all tested case are converged
- better cache usage (gain ~ 2)
- less iteration (gain ~ 3)

# Summary

- introduced QCD, LQCD and Krylov subspace methods briefly
- Modified Block BiCGSTAB showed remarkable cost reduction

# Summary

- introduced QCD, LQCD and Krylov subspace methods briefly
- Modified Block BiCGSTAB showed remarkable cost reduction
- and should accelerate LQCD simulations on

